## Neutron Star Merger Simulation - our latest progress -



Masaru Shibata



Max Planck Institute for Gravitational Physics at Potsdam & Yukawa Institute for Theoretical Physics, Kyoto U.

with the help of S. Fujibayashi, K. Hayashi, K. Kiuchi & S. Wanajo 23.06.2022 binary neutron stars workshop

### Outline

- I. Introduction
- II. NS-merger & post-merger simulations separately
- III. Seconds-long BHNS-merger simulations IV. Issues

### I Introduction

Neutron star mergers are predicted to be

- Promising source for short-hard gamma-ray bursts (e.g., Eichler et al. 1989)
- Site for *r*-process nucleosynthesis (rapid neutron capture nucleosynthesis) (e.g., Schramm & Lattimer 1974)
- Source for Kilonovae/macronovae (e.g., Li & Paczynskii, 1998)
- Invaluable site for studying nuclear equation of state through GW detection (e.g., Lai et al. '93, Hinderler, '08,....)
- GW170817 (1<sup>st</sup> NS-NS) has shown all these aspects



#### Variety of NS-NS merger process





Sneden et al. (2008):

**Can NS merger reproduce this?** 

Courtesy of Sekiguchi's presentation

#### Variety of NS-NS merger process



#### **Roles of merger simulations in NR**

- 1. To provide information to predict signals (EM, GW, CR,  $\nu$ ) associated with mergers for the future events
- 2. To explore nucleosynthesis: e.g, whether the universality is reproduced & what are observable lines for heavy elements? (talk by Hotokezaka)
- To prepare a setup for the post-process simulations;
  e.g., long-term evolution of ejecta associated with kilonova & its afterglow, jet simulations, etc.

#### Time line after the merger: *to be studied* $\gamma$ -ray burst? (< 2s) compact obj + disk Merger $\rightarrow$ 100 10 1000 ms Time after merger Dynamical ejection by shock heating/tidal effect (proceeds in < 10 ms) **Key: Angular momentum transport** MHD/viscosity-driven post-merger ejection (in MHD/viscous timescale of remnant disk $\sim 1$ s)

Key: Weak interaction which determines the property of ejecta: Important for nucleosynthesis & kilonovae (macronovae)

> Electron fraction,  $Y_e(=n_p/(n_n+n_p))$ , is key quantity for r-proc.

#### A milestone goal of NS merger simulations

- Important timescales:
- 1. Dynamical mass ejection timescale  $\sim 10$  ms
- 2. Post-merger mass ejection timescale ~ O(1) s
- 3. Short gamma-ray bursts: typically 1~2 s
- → We need to evolve the post-merger system at least for ~10 sec
- Input physics is also important
- *MHD/viscous effects* are the key to PM evolution
- Weak processes are the key to nucleosynthesis

→ Seconds-long GR+rad+MHD simulation with weak physics input (e.g., neutrino transfer) is needed

#### **Two ways**

- Merger (<~100 ms) + post-merger simulations separately (e.g., talk by Francois)
- 2. Seconds-long self-consistent simulation

✓ Nucleosynthesis, kilonova & its afterglow, GRB jets are analyzed as a post-process

#### **II Merger & post-merger separately**

- We have been working on this for the last five years: Popular way currently
- A solid way to specifically explore two mass ejection mechanisms

## A Dynamical mass ejection from NS-NS (ejection within ~ 10 ms after the merger)

- Many *numerical-relativity* simulations have been done since 2013 → Well understood besides quantitative details
- ♦ What we have learned are
- Mass= $10^{-4} \sim 10^{-2} M_{sun}$  (Hotokezaka, Sekiguchi, Foucart, Radice...., now it is routine work): For *low total mass (MNS formation)*, it is  $<\sim 10^{-3} M_{sun}$  but could be higher for HMNS case
- Electron fraction=0.05~0.4 (show later)
  → suitable for *r*-process nucleosynthesis of heavy elements (Wanajo, Sekiguchi, Goriely, Foucart, Roberts, and others)
- Average velocity=0.15~0.25c, but could be up to
  ~ 0.9c (or more) (Hotokezaka+ '13, many follow-ups, Radice....)

#### **B** Post-merger mass ejection: *more complicated*

- Neutron star is *magnetized* → Remnants are magnetized
- The magnetic field is amplified by **MHD instabilities** (Kelvin-Helmholtz instability, MRI, convection, etc)
- i. Turbulence & effective viscosity are excited (Fernandez & Metzger+ '13, Just et al. '15, '21, Fujibayashi+ '18, '20)
- ii. Purely MHD effects (e.g., Christie+ '19, Just+ '21, Shibata+ '21)
  → Post-merger mass ejection from disk/torus
- ✓ Ejecta mass depends on the remnant (BH or NS) →  $M_{eje} \sim 0.05-0.1 M_{sun}$  for long-lived NS formation, while it is lower, ~ 0.01 M<sub>sun</sub>, for BH formation
- Weak interaction physics (e.g., neutrino reaction) is key for determining electron fraction (Y<sub>e</sub>) (e.g., Metzger & Fernandez '14, Just+ '15, '21, Fujibayashi+ '18, '20, Miller '19)

# Basic evolution process of disks by neutrino cooling and (effective) viscous effects

 $T_{max} > 3 \text{MeV} \rightarrow L_{\nu} \approx \dot{E}_{vis}$ No viscous mass ejection; Viscous angular momentum transport  $\rightarrow$  Disk expansion (but no mass ejection)

BH/NS

*ک*ر کر ۷

 $T_{max} < 3 \text{MeV} \rightarrow L_{\nu} \ll \dot{E}_{vis}$ Viscous heating is fully used for matter expansion  $\rightarrow$ Onset of viscous mass ejection

<u>Viscous angular momentum transport timescale ~ sec</u>  $\tau_{\rm vis} \sim 0.55 \,\mathrm{s} \left(\frac{\alpha_{\rm vis}}{0.02}\right)^{-1} \left(\frac{\varpi}{50 \,\mathrm{km}}\right)^2 \left(\frac{c_s}{0.05c}\right)^{-1} \left(\frac{H}{20 \,\mathrm{km}}\right)^{-1}$ 

#### Viscous hydro simulation in full GR: 3 solar mass BH + 0.1 solar mass disk ( $\alpha$ =0.05)



#### How $Y_e$ of disks/ejecta is determined?

- <u>β-equilibrium</u> (reaction timescale < disk evolution one)</li>
  p + v<sub>e</sub> ↔ n + e<sup>+</sup> & n + v<sub>e</sub> ↔ p + e<sup>-</sup>
  → Y<sub>e</sub> is determined by μ<sub>p</sub> + μ<sub>e</sub> = μ<sub>n</sub> + μ<sub>y</sub>
- In typical situations, neutrino captures decouple first, but still **electron & positron capture processes proceed** because of high temperature > MeV

 $p + e^- \rightarrow n + \nu_e \& n + e^+ \rightarrow p + \overline{\nu_e}$ 

- For  $T_{\text{max}} < \sim 3$ MeV, the weak interaction decouples and  $Y_{\text{e}}$  is determined (Fujibayashi+ '20, Just+ '21)
- ✓ Electron degeneracy is weakened for decreased density, i.e.,  $\mu_e$  decreases with time → At mass ejection, moderately neutron rich,  $Y_e \sim 0.3$ → Heavy r elements production is suppressed
  - $\rightarrow$  Heavy *r*-elements production is suppressed



Numerical relativity results by Fujibayashi et al. arXiv: 2205.05557

Mass ratios by *dynamical* and *post-merger ejecta* depend significantly on the lifetime of remnant NS

#### **Broadly speaking there are four patterns** $dM/dY_{e}$ $dM/dY_{e}$ 0.2-0.3 Disk mass 0.2-0.3 is $\sim 0.2$ solar Disk mass High PM ejecta mass is ~0.05 solar Low dynamical ejecta mass $Y_{\rm e}$ $Y_{\rm e}$ 0.4 0.4 <u>NS-NS $\rightarrow$ long-lived NS (Magnetar)</u> Symmetric NS-NS $\rightarrow$ short-lived NS $dM/dY_{\rm e}$ $dM/dY_{e}$ 0.2-0.3 0.2-0.3 Disk mass is Disk mass 3—7 dynamical is $>\sim 0.1$ solar ejecta mass No shock heating Less shock heating $Y_{\rm e}$ $Y_{\rm e}$ 0 0.4 0.4 Asymmetric NS-NS→ short-lived NS $BH-NS \rightarrow BH + disk$

#### **Nucleosynthesis for MNS vs BH formation** See Fujibayashi et al. 2022; arXiv: 2205.05557





## MHD effect does not significantly change the abundance pattern for long-lived MNS formation case



Shibata et al. PRD 104, 063026 (2021); Kawaguchi et al. ApJ in press

#### Prediction

- Only when the remnants of NS-NS mergers collapse to a BH in a short timescale (~10—100ms) after the merger, the *universality* can be reproduced
- Suppose that the NS-NS merger is the main site for the *r*-process nucleosynthesis. Then,
   <u>long-lived MNS formation should be a minority</u>
- → Short-GRBs from a remnant magnetar should be rare; consistent with the absence of radio sources (talk of KH)

#### In other words,

• In case many long-lived MNS formations are observed in the near future, the NS-NS merger scenario of the *r*process nucleosynthesis may be excluded.

See Fujibayashi et al. 2022; arXiv: 2205.05557

#### **III Seconds-long BH-NS merger simulation**

K. Hayashi et al. arXiv: 2111.04621; PRD in press

• With the current computational resources, selfconsistent simulations are feasible at least for BH-NS binaries because the length scale is larger than the NS-NS case:

> $\Delta t \propto \Delta x \propto M_{BH} (< \sim 10 M_{\odot})$ Focusing only on the case that NS is tidally disrupted

#### BH-NS merger for 2 seconds: GR + v-rad + MHD

NS with strong dipole field initially

K. Hayashi et al. arXiv: 2111.04621



y=0 plane is displayed: [-2000,2000]km  $\Delta x=400m$ ; Fix mesh refinement with ~400\*400\*200 grid \* 9 levels







#### 3D view: By Kota Hayashi

*NEW*: Absence of equatorial-plane symmetry



Rest-mass density + magnetic field lines

#### Magnetic-field strength

### **IV** Issues

- Next step: Perform seconds-long simulations for binary neutron stars: Kiuchi is working on this now
- Improve weak interaction
- ✓ Better neutrino radiation transfer (never-ending issue)
- ✓ Neutrino oscillation (only qualitative study possible)
- Improve angular momentum transport
- ✓ MHD with high resolution (never-ending issue)
- <u>Dynamo in remnant neutron stars</u> (long-term simulation is needed; not yet seriously attacked)
- ✓ How global magnetic fields are developed in the presence of remnant neutron stars?

#### **Electron fraction in resistive MHD+ dynamo terms**



Shibata et al. PRD 104, 063026 (2021); Kawaguchi et al. ApJ in press